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# Gaseous CO<sub>2</sub> Behaviour during Water Displacement in a Sandstone Core Sample

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**Abstract:** CO<sub>2</sub> injection into subsurface formations involves the flow of CO<sub>2</sub> through a porous medium that also contains water. The injection, displacement, migration, storage capacity and security of CO<sub>2</sub> is controlled mainly by the interfacial interactions and capillary, viscous, and buoyancy forces which are directly influenced by changes in subsurface conditions of pressure and temperature; the impact of bouncy forces is assumed negligible during this study. In this study, gaseous CO<sub>2</sub> is injected into a water-saturated sandstone core sample to explore the impact of fluid pressure (40-70 bar), temperature (29-45 °C), and CO<sub>2</sub> injection rate (0.1-2 ml/min) on the dynamic pressure evolution and displacement efficiency. This study highlights the impact of capillary or viscous forces on the two-phase flow characteristics and shows the conditions where capillary or viscous forces become more influential. The results reveal a moderate to considerable impact of the parameters investigated on the differential pressure profile, endpoint CO<sub>2</sub> relative permeability ( $K_{rCO_2}^{max}$ ), and irreducible water saturation ( $S_{wr}$ ). Overall, the increase in fluid pressure, temperature, and CO<sub>2</sub> injection rate cause an increase in the maximum and final differential pressures, an increase in the  $K_{rCO_2}^{max}$ , a reduction in the  $S_{wr}$ .  $S_{wr}$  was in the range of around 0.38-0.45 while  $K_{rCO_2}^{max}$  was less than 0.25. The data show a significant influence for the capillary forces on the pressure and production behaviour. The capillary forces produce high oscillations in the pressure and production data while the increase in viscous forces impedes the appearance of these oscillations. The appearance and frequency of the oscillations depend on the fluid pressure, temperature, and CO<sub>2</sub> injection rate but to different extents.

# 1 Introduction

Carbon capture and storage (CCS) is regarded as one of the most promising techniques that can deal effectively with the increasing emissions of anthropogenic CO<sub>2</sub> into the atmosphere due to fossil fuel burning and other human activities (Bachu, 2001; Hangx et al., 2013; Kazemifar et al., 2015). The captured CO<sub>2</sub> can be sequestered in deep saline aquifers, depleted or abandoned oil and gas reservoirs (Delshad et al., 2010; Gozalpour et al., 2005; Kaveh et al., 2012), or unmineable coal bed seams (Kaveh et al., 2012; Plug and Bruining, 2007) to enhance recovery from hydrocarbon reservoirs, increase methane production from coal beds, or extract geothermal heat from subsurface formations (Kaveh et al., 2012; Tutolo et al., 2015). Figure 1 presents a summary of the pressure and temperature ranges at which saline aquifers are found underground and highlights that CO<sub>2</sub> can exist in a gaseous, liquid or supercritical phase (Bachu, 2000; Espinoza and Santamarina, 2010; Frailey et al.; Nourpour Aghbash and Ahmadi; Saraji et al., 2014; Sohrabi et al.).

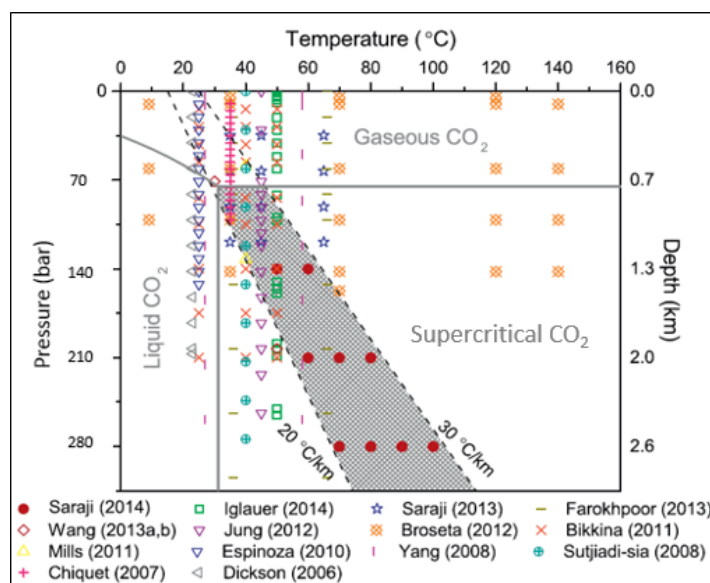


Figure 1: The pressure and temperature ranges at which saline aquifers are found underground (Saraji et al., 2014). This study is conducted under pressure ranged from 40 to 70 bar and temperature ranged from 29 to 45 °C.

During CO<sub>2</sub> injection in subsurface formations, the bulk of the injected CO<sub>2</sub> (as a non-wetting fluid) will displace the formation water (as a wetting fluid) in an immiscible displacement (Basbug et al., 2005; Herring et al., 2014b). The displacement of the injected CO<sub>2</sub> depends on a number of parameters,

namely, the interfacial interactions (e.g. interfacial tension and wettability), solubility of CO<sub>2</sub> in formation water, densities and viscosities of fluids present, petrophysical properties of the subsurface formation, injection rate and its duration , and more importantly on the capillary and viscous forces ([Cinar and Riaz, 2014](#); [Duan and Sun, 2003](#); [Pentland et al., 2011b](#); [Trevisan et al., 2017](#)). The capillary forces at the CO<sub>2</sub>-water interface are of considerable importance in determining the nature of the flow through pores ([Roof, 1970](#)). Any change in subsurface conditions of pressure and temperature will have a significant impact on the interfacial interactions ([Espinoza and Santamarina, 2010](#); [Liu et al.](#); [Plug and Bruining, 2007](#); [Yang et al., 2007](#)), the viscous forces due to the change in viscosity ([Bachu and Bennion, 2008b](#)) and the capillary forces. The change in interfacial interactions, and viscous and capillary forces due to the change in underground conditions will have a considerable influence on the capillary pressure, relative permeability ([Alkan et al., 2010](#)), pore-scale fluid distribution ([Al-Menhali and Krevor, 2014](#)), CO<sub>2</sub> injection, fluid migration, capacity and long-term fate of CO<sub>2</sub> storage in saline aquifers ([Levine et al., 2011](#); [Saraji et al., 2013](#); [Wang et al., 2015](#)), CO<sub>2</sub>-enhanced oil and gas recovery processes ([Gozalpour et al., 2005](#); [Qi et al., 2010](#)). According to Salimi et al., the change in capillary pressure, due to the change in the operational conditions, can have a direct influence the CO<sub>2</sub>-storage capacity and the heat recovery due to its impact on the solubility and density of both CO<sub>2</sub> and water ([Salimi et al., 2012](#)). Thus, it is of utmost importance to have a deep insight into the dynamic behaviour of CO<sub>2</sub> under different operational conditions.

CO<sub>2</sub> has been used in the oil industry for a long time, in particular, to increase productivity through Enhanced Oil Recovery (EOR), and extensive research has been undertaken describing multi-phase flow properties of CO<sub>2</sub>-oil systems ([Bahralolom et al., 1988](#)). On the other hand, much less laboratory investigations have been done for CO<sub>2</sub>-water (brine) systems ([Perrin and Benson, 2010](#)). Those published have mainly focused on CO<sub>2</sub> wettability ([Al-Menhali and Krevor, 2014](#); [Bikina, 2011](#); [Farokhpour et al., 2013a](#); [Kaveh et al., 2012](#); [Sakurovs and Lavrencic, 2011](#); [Saraji et al., 2013](#)), CO<sub>2</sub>-water (brine) interfacial tension ([Aggelopoulos et al., 2010](#); [Bachu and Bennion, 2008b, 2009](#); [Busch and](#)

[Müller, 2011](#); [Chiquet et al., 2007](#); [Li et al., 2012](#); [Yu et al., 2012](#)), relative permeability ([Bachu, 2013](#); [Krevor et al., 2015](#); [Liu et al.](#); [Perrin et al., 2009](#)) and capillary pressure ([Busch and Müller, 2011](#); [Pini et al., 2012](#); [Plug and Bruining, 2007](#)). Cinar and Riaz showed that much of the research has been directed to investigate the fluid properties rather than studying the multiphase flow properties of the CO<sub>2</sub>-water systems ([Cinar and Riaz, 2014](#)).

The limited investigations of the multiphase flow characteristics of CO<sub>2</sub>-water (brine) systems have involved laboratory experiments ([Jobard et al., 2013](#)), computational modelling ([Jobard et al., 2013](#); [Ma et al., 2013](#); [Xu et al., 2011](#)), and field scale projects ([Wang et al., 2015](#)). The CO<sub>2</sub>-water (brine) investigations included core flooding displacements performed at liquid, supercritical and gaseous CO<sub>2</sub> conditions. Current literature survey of the CO<sub>2</sub>-water (brine) multiphase flow experiments showed that most of these experiments were supercritical (Sc) CO<sub>2</sub>-brine (water) displacements studies, which were performed on various porous media such as core samples ([Berg et al., 2013](#); [Edlmann et al., 2013](#)), micromodels ([Cao et al., 2016](#)), and packed beds of glass beads ([Song et al., 2012](#); [Suekane et al., 2005](#)). In these studies related to supercritical CO<sub>2</sub> migration, researchers have examined various parameters such as relative permeability curves ([Berg et al., 2013](#); [Chang et al., 2013](#); [Krevor et al., 2013](#); [Suekane et al., 2005](#); [Suenaga and Nakagawa, 2011](#)), capillary pressure curves ([Herring et al., 2014a](#); [Wang et al., 2013](#)), CO<sub>2</sub> residual saturation and distribution ([Alemu et al., 2011](#); [Chang et al., 2013](#); [Herring et al., 2014a](#); [Pentland et al., 2011a](#); [Saeedi et al., 2011](#); [Suekane et al., 2005](#)), heterogeneity impact ([Ott et al., 2015](#); [Perrin and Benson, 2010](#); [Shi et al., 2011](#); [Wang et al., 2013](#)), water displacement efficiency ([Cao et al., 2016](#)), mass transfer ([Berg et al., 2013](#)), and formation dry-out ([Ott et al., 2011](#)). Some liquid (L) CO<sub>2</sub>-water (brine) core flooding displacements were conducted to investigate the multiphase flow characteristics of CO<sub>2</sub>-water-porous media ([Manceau et al., 2015](#)), CO<sub>2</sub> residual saturation and distribution ([Alemu et al., 2011](#)), and pore-scale heterogeneity ([Zhang et al., 2011](#)).

On the other hand, very scarce data was found regarding gaseous (G) CO<sub>2</sub> injection into water (brine) saturated porous systems ([Islam et al., 2013](#); [Jiang et al., 2017](#); [Lassen et al., 2015](#); [Yu et al., 2014](#)).

Even though liquid and supercritical CO<sub>2</sub> injection is more efficient, the dynamic behaviour of gaseous CO<sub>2</sub> in reservoir rock is necessary information, particularly considering that many potential saline storage aquifers are within temperature and pressure conditions of the gaseous CO<sub>2</sub> phase (Figure 1) and that any leakage of CO<sub>2</sub> from deeper storage would inevitably result in a phase change to a gaseous CO<sub>2</sub> state (Edlmann et al., 2016; Miocic et al., 2016). The existing GCO<sub>2</sub>-water experiments were designed to investigate the crossover zone of flow regimes, impact of capillary number, CO<sub>2</sub> injection rates and permeability on displacement efficiency. Islam et al. conducted GCO<sub>2</sub>-water experiments at 1 bar and 25 °C using a vertical Hele-Shaw cell filled with micro-beads to investigate the crossover zone from capillary to viscous to fracture fingering. They observed that all the three fingering patterns can occur in the cell but at different heights (Islam et al., 2013). Jiang et al. performed both immiscible and miscible drainage GCO<sub>2</sub>-water displacements inside a packed bed filled with quartz glass beads to have a better understanding of the two-phase flow characteristics inside porous media. The experiments were conducted at CO<sub>2</sub> injection rates varying from 0.01 to 3 ml/min and at 60 bar and 24.85 °C. They observed that: (I) at low CO<sub>2</sub> injection rates, the CO<sub>2</sub> dissolution increases; (II) the increase in glass beads diameter (i.e. higher permeability) leads to a decrease in the capillary forces (Jiang et al., 2017). Yu et al. conducted immiscible drainage GCO<sub>2</sub>-water displacements at 60 bar and 24.85 °C inside a packed bed of glass beads (0.2 mm diameter) to study the impact of the capillary number on displacement efficiency. They noticed that the increase in the capillary number, when it is between 10<sup>-11</sup> and 10<sup>-10</sup>, results in a sharp reduction in the residual water saturation as a result of increasing the impact of the viscous forces (Yu et al., 2014).

Despite the considerable research on the CO<sub>2</sub>-water (brine) systems and its practical importance, the analysis of the pressure data in core flooding has been widely overlooked (Rezaei and Firoozabadi, 2014). To the authors' best knowledge, there is no detailed investigation into the dynamic pressure evolution and displacement efficiency of gaseous CO<sub>2</sub> during its injection into a water saturated core sample. In this paper, laboratory dynamic drainage experiments were performed by injecting pure CO<sub>2</sub> into the deionised water-saturated sandstone core sample to investigate the impact of fluid pressure,

temperature, and CO<sub>2</sub> injection rate on the differential pressure profile, water production, and endpoint effective and relative permeabilities of CO<sub>2</sub>. This study also highlights the impact of capillary and viscous forces on the pressure and production data as well as shows the conditions at which capillary or viscous forces become more influential. During these dynamic displacements, the transient pressure at the inlet and outlet sides of the core and the transient outflow rates of water and CO<sub>2</sub> were measured and analyzed. The endpoint water saturations of CO<sub>2</sub> and water were also calculated.

## 2 Materials

A sandstone core sample from the Guillemot A Field in the North Sea was used to perform the unsteady state GCO<sub>2</sub>-water drainage experiments. The core sample has a diameter of 2.54 cm and a length of 7.62 cm. The average porosity and absolute water permeability of the core sample were about 14% and 15.8 millidarcys, respectively. This study is one in a series, thus the core sample description, the experimental setup and the CO<sub>2</sub>-water displacement procedures can be seen in our recent publication ([Al-Zaidi et al., 2018](#)).

## 3 Results and discussion

To gain a deep insight into the dynamic behaviour of GCO<sub>2</sub>-water drainage displacements under various fluid pressure, temperature, and injection rate conditions; the inlet and outlet pressure, CO<sub>2</sub> and water out flowrate, the irreducible water saturation and endpoint effective and relative permeabilities of CO<sub>2</sub> were measured and analyzed.

In this study, the difference between the pressure transducer readings at the inlet and outlet sides of the core sample has been used to calculate the differential pressure. The differential pressure during horizontal CO<sub>2</sub> injection is largely influenced by the capillary and viscous forces. The capillary forces are controlled mainly by the CO<sub>2</sub>-water interfacial tension, contact angle (i.e. wetting status), pore diameter and geometry ([Alkan et al., 2010](#); [Bikkina et al., 2016](#); [Chatzis and Morrow, 1984](#); [Fulcher Jr et al., 1985](#)). The wetting status plays an important role in determining the imbibition and the distribution of the wetting and non-wetting phases inside the porous media ([Chalbaud et al., 2007](#); [Espinoza and](#)

[Santamarina, 2010](#)). The capillary forces, which are responsible for the entrapment of one phase by another during immiscible displacements in porous media ([Akbarabadi and Piri, 2013](#); [Chatzis and Morrow, 1984](#)), arise from the presence of the interface between the immiscible fluids ([Bikkina et al., 2016](#)) and significantly dominate the multiphase flow, especially in low permeability rocks and fractured reservoirs ([Schembre and Kovscek, 2003](#)). On the other hand, the viscous forces are controlled mainly by the viscosity of both displacing and displaced fluids, the fluid velocity in the pores, the amount of each fluid (i.e. saturation) in the pore, and the core sample properties (e.g. frontal area, permeability, and length). Espinoza and Santamarina ([Espinoza and Santamarina, 2010](#)) proposed the following equation to account for the impact of the capillary and viscous forces on the differential pressure as follow:

$$\Delta P = P_{CO_2} - P_{water} = 4 \frac{\sigma_{CO_2-water} \cos \theta}{d} + v \frac{32 L}{d^2} \left( \frac{l_{CO_2} \mu_{CO_2} + l_{water} \mu_{water}}{L} \right) \quad (1)$$

Where  $\Delta P$  is the differential pressure across the core sample (Pa).  $P_{CO_2}$  and  $P_{water}$  are the pressures of CO<sub>2</sub> phase and water phase, respectively.  $\sigma_{CO_2-water}$  is the CO<sub>2</sub>-water interfacial tension (mN/m),  $\theta$  the contact angle,  $d$  (m) the diameter of the largest effective pore ([Chiquet et al.](#); [Chiquet et al., 2007](#); [Farokhpour et al., 2013b](#); [Han et al., 2010](#)),  $L$  (m) the length of the core sample,  $l$  (m) the length of CO<sub>2</sub> or water phase inside the core sample,  $v$  (m/s) the fluid velocity in the pores, and  $\mu$  (Pa.s) the viscosity of the fluids. The first term of Eq.1 refers to the Young-Laplace equation, which accounts for the capillary forces, while the second term refers to the Poiseuille's equation ([Espinoza and Santamarina, 2010](#); [Li, 2015](#)), which account for the viscous forces. For small injection rate and high viscosity contrast conditions [the impact of viscous forces can be neglected](#), thus Eq.1 can be reduced to the Young-Laplace equation ([Li, 2015](#)) as follows:

$$\Delta P = P_{CO_2} - P_{water} = 4 \frac{\sigma_{CO_2-water} \cos \theta}{d} \quad (2)$$

The Young-Laplace equation is used to determine the critical pressure point, which is the differential pressure required for the displacing fluid to enter the core sample for the first time. The



non-wetting fluid cannot enter the core sample unless its pressure becomes higher than the critical pressure point ([Han et al., 2010](#)).

In this study, the experimental results have been categorized into two main sections. The first section presents and discusses the impact of the experimental fluid pressure, temperature and CO<sub>2</sub> injection rate on the differential pressure profiles while the second section deals with the impact of the parameters investigated on the endpoint CO<sub>2</sub> effective (relative) permeability and irreducible water saturation.

It should be noted that during this study, the term low and high-fluid pressure refers to the experiments conducted at pressures less and higher than 50 bar, respectively. The low and high temperature refers to the experiments performed at less or higher than 33 °C, respectively. The low, medium and high injection rates refer to the experiments performed at injection rate ranging from 0.1 to 0.2 ml/min, from 0.3 to 0.6 ml/min, and from 1 to 2 ml/min, in sequence. The corresponding time refers to the time required to reach the maximum-differential pressure at the start of the experiment. The quasi-differential pressure refers to the differential pressure at the end of the experiment.

### ***3.1 Differential Pressure Profile of GCO<sub>2</sub>-Water Drainage Displacements***

To investigate the effect of fluid pressure, experimental temperature, and CO<sub>2</sub> injection rate on the differential pressures, series of GCO<sub>2</sub>-water displacements were performed at various fluid pressures (from 40 to 70 bar), experimental temperatures (29-45 °C) and CO<sub>2</sub> injection rates (0.1-2 ml/min).

#### ***3.1.1 Effect of Fluid Pressure on the Differential Pressure Profile of GCO<sub>2</sub>-Water Drainage Displacements***

[Figure 2](#) presents the impact of increasing fluid pressure on the differential pressure profile of GCO<sub>2</sub>-water drainage displacements. A number of trends are identifiable: Firstly, the differential pressure profile at all fluid pressures is characterized by a high initial increase, immediately followed by a steep rapid reduction and then followed by a quasi-differential pressure. Secondly, there are

multiple oscillations of these cycles. The frequency of these oscillating cycles increases as fluid pressure increases along with a rise in the values of the maximum and quasi-differential pressures.

The high initial increase in the differential pressure can be related to the capillary pressure. The following reduction in the differential pressure profile reflects the impact of the reduction in both capillary forces and viscous forces according to Eq.1. The injection of gaseous CO<sub>2</sub> into the core sample generates the initial increase in differential pressure to overcome the capillary entry pressure for the invasion of the gaseous CO<sub>2</sub> ([Chang et al., 2013](#)). The reduction in the capillary forces can be associated with the reduction in the pore resistance to CO<sub>2</sub>-water interfaces as the number of pores opened by CO<sub>2</sub> is increased ([Kwelle, 2017](#)). This agrees very well with Kwell's finding, who noticed a high reduction in the differential pressure profile as the CO<sub>2</sub>-water interface is displaced out of microcapillary tubes ([Kwelle, 2017](#)). The reduction in the viscous forces can be related to the combined effect of the dynamic change in relative permeability of gaseous CO<sub>2</sub> and water and the high rate replacement of a more viscous fluid (water) with a less viscous fluid (CO<sub>2</sub>) ([Chang et al., 2013](#)). [Replacing water by CO<sub>2</sub> at a high rate](#) can be linked to (a) the high mobility ratio due to the high viscosity contrast and (b) gas expansion effects which generate an increase in volumetric CO<sub>2</sub> injection rate inside the core sample.

- The gas expansion can, in turn, be related to the density change of the injected CO<sub>2</sub> due to the temperature difference between inside the water bath (i.e. 29 to 45 °C depending on the experimental conditions) and outside it (room temperature 18-20 °C). The density of the injected CO<sub>2</sub> varies as the CO<sub>2</sub> enters the water bath dependant on the injection rate, fluid pressure and the temperature difference from the pump to the sample. The density ratio ( $d_r$ ) suggested by Perrin and Benson ([Perrin and Benson, 2010](#)) has been used to calculate the injection rate inside the core sample. For instance, at an experimental pressure of 40 bar, an injection rate of 1 cm<sup>3</sup>/min at 20 °C becomes 1.7522 cm<sup>3</sup>/min at 33 °C. However, at an experimental pressure of 70 bar and the same injection rate and temperature conditions, it becomes 5.281cm<sup>3</sup>/min.

$$d_r = \frac{P_1 T_2 Z_2}{P_2 T_1 Z_1} \quad (3)$$

Figure 2 reveals that the differential pressure profiles are characterized by multiple differential pressure (PD) oscillations. The appearance of these PD oscillations can be related to the impact of the capillary forces at the trailing end of each CO<sub>2</sub>-water slug during CO<sub>2</sub> flooding (Nutt, 1982) or the capillary end effects. According to Nutt, the impact of the capillary forces at the trailing end of the CO<sub>2</sub>-water slug is governed by the wetting status of the injected fluid. If a non-wetting fluid (e.g. CO<sub>2</sub>) is injected, then the capillary forces will work in an opposite direction to the applied viscous forces. Thus, as water depletion is progressed, the applied viscous forces will drop until they become less than the capillary forces. Upon reaching this point, the flow of the non-depleted capillaries is potentially blocked by the capillary forces (Nutt, 1982). This blockage occurs due to a re-imbibition process of the wetting phase inside the core sample, which was noticed by Hildenbrand et al (Hildenbrand et al., 2002). Hildenbrand et al. observed that the re-imbibition process occurs when the excess pressure in the non-wetting phase declines after the gas breakthrough (Hildenbrand et al., 2002), as shown in Figure 3. This re-imbibition process occurs in a progressive manner starting with the smallest pores and continuing to the larger pores, leading to the successive loss of the interconnected flow-paths, which, in turn, leads to a progressive decline in the non-wetting phase relative permeability. Finally, when the last interconnected flow-path for the non-wetting phase is blocked, the permeability of the non-wetting phase will drop to zero (Hildenbrand et al., 2002). According to Hildenbrand et al., this re-imbition process can result in a residual water saturation when certain-gas filled pores become isolated a result of interrupting the flow pathways. The maximum differential pressure required to open the flow paths again can be used to determine the largest effective pore radius and, hence, the sealing efficiency of the rock (Hildenbrand et al., 2002).

Therefore, since our core sample is water-wet, the pressure of the injected CO<sub>2</sub> had to build up to a certain level to overcome the capillary forces that blocked the CO<sub>2</sub> outflow rate (Nutt, 1982). Due to the high compressibility nature of the gaseous CO<sub>2</sub>, the injected CO<sub>2</sub> will accumulate inside the core

sample and the connections pipes until the differential pressure becomes high enough to overcome the capillary forces. Once the blocked capillaries are opened to flow, the cumulative CO<sub>2</sub> will expel the liquid drops that block the pores out of the core sample quickly; the rate of [expulsion](#) is expected to increase with the fluid pressure. The development of this phenomenon is highly influenced by the core sample properties and the injection rate due to their direct impact on viscous and capillary forces. As a result, this phenomenon is expected to be reduced when the injection rate, i.e. viscous pressure drop, becomes high enough to overcome the capillary forces ([Nutt, 1982](#)). However, due to the cyclic reduction of the viscous pressure drop (i.e. viscous forces) to the level that becomes insufficient to overcome the capillary forces, this phenomenon of oscillations can occur frequently.

On the other hand, since the GCO<sub>2</sub>-water displacements are strongly influenced by the capillary end effects and viscous instabilities ([Müller, 2011](#)), it might be suggested that the appearance of the oscillations is due to the impact of capillary end effects. The capillary end effects occur at both inlet and outlet faces of the core sample, but its impact becomes more severe at the outlet face. According to Müller, the capillary end effects can never be entirely prevented but can be corrected for ([Müller, 2011](#)). The impact of capillary end effects and viscous instabilities can be reduced when the following scaling coefficient proposed by Rapoport and Leas for stabilized floods becomes greater than one.

$$Lu\mu \geq 1 \quad (4)$$

where  $L$  is the length of the medium (cm),  $u$  the Darcy velocity (cm/min), and  $\mu$  the displacing phase viscosity (cp) ([Fathollahi and Rostami, 2015](#)). The scaling coefficients for the 40, 50, and 70 bar displacements are 0.0773, 0.0844, and 0.285, respectively. The scaling coefficients increased significantly as the fluid pressure increased from 40 and 50 bar to 70 bar, which indicates a reduction in the impact of capillary end effects with increasing fluid pressure. However, since the data from [Figure 2](#) reveal an increase in the frequency of the oscillations with increasing fluid pressure, this indicates that the capillary end effects are not responsible for the PD oscillation phenomenon. In addition, the

disappearance of the oscillations at lower injection rates as shown in [Figure 7](#) further supports the idea that the oscillations are not because of the capillary end effects.

[Figure 2](#) also shows that increasing fluid pressure leads to an increase in the rate of the differential pressure (PD) oscillations along with increases in the values of the maximum and quasi-differential pressures and a reduction in the corresponding time (the time required to reach the maximum-differential pressure at the start of the experiment). For illustration, it can be seen that as the fluid pressure increased from 40 to 50 bar, the rate of the PD oscillations increased by around 33% and the maximum-differential pressure increased by about 2.50%. The quasi-differential pressure was constant at around 1 bar. The corresponding time declined by approximately 17%. However, as the fluid pressure increased from 50 to 70 bar, the PD oscillations substantially increased by 225%, the maximum-differential pressure raised by around 9% and the quasi-differential pressure increased by 165%. The corresponding time dropped considerably by around 78%. The high reduction in the corresponding time with increasing fluid pressure can be related mainly to the increase in gaseous CO<sub>2</sub> density and the injection rate inside the core sample due to the expansion effects. [As gaseous CO<sub>2</sub> becomes denser, it requires lesser time to be compressed to the required pressure.](#)

The increase in the maximum and quasi-differential pressures with increasing fluid pressure can be related mainly to the magnitudes of both viscous and capillary forces. [According to Eq.1, as the fluid pressure increases the viscous forces increase \[due to the increase in CO<sub>2</sub> viscosity and the injection rate inside the core sample due to expansion impact\], while the capillary forces decrease \[because of the reduction in the CO<sub>2</sub>-water interfacial tension \(IFT\) \(\[Georgiadis et al., 2010\]\(#\)\) and the increase in the contact angle \(\[Banerjee et al., 2013\]\(#\)\) due to increasing CO<sub>2</sub> solubility \(\[Bennion and Bachu; Yang et al., 2007\]\(#\)\)\]](#). Thus, the increase observed in the differential pressures is the net result of the increase in the viscous forces and the reduction in the capillary forces. Reducing capillary forces with increasing pressure is expected to cause a reduction in the extent of differential pressure increase.

The increase in the PD oscillations means the frequency of liquid drops expelled out of the core sample is increased. This can be associated mainly with the reduction in the capillary forces and the increase in gas density with increasing pressure. Increasing the gas density and reducing capillary forces mean less time was required to reach a differential pressure value which was sufficient to overcome the capillary forces; thus, increasing the frequency of the PD oscillations.

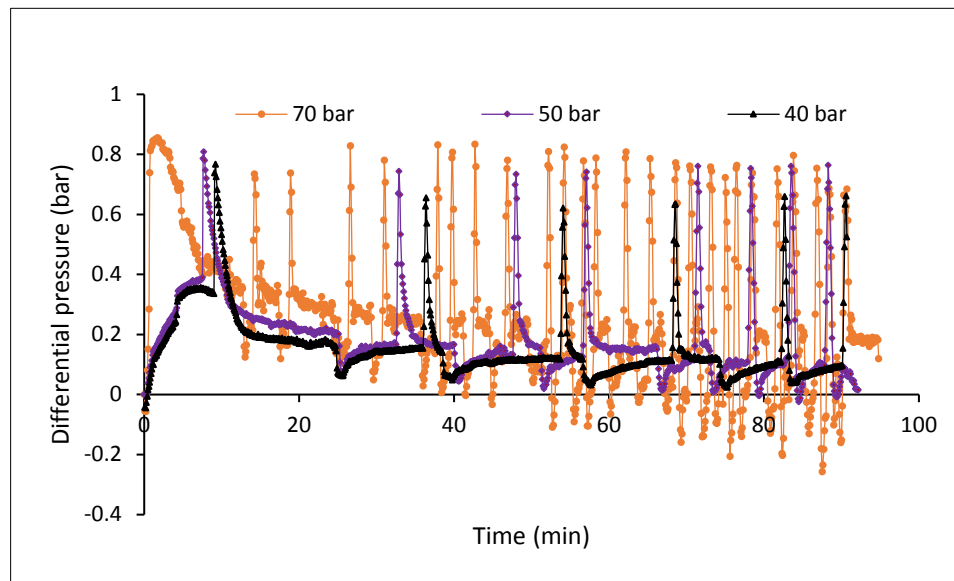


Figure 2: Effect of fluid pressure on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 0.4 ml/min and 33 °C.

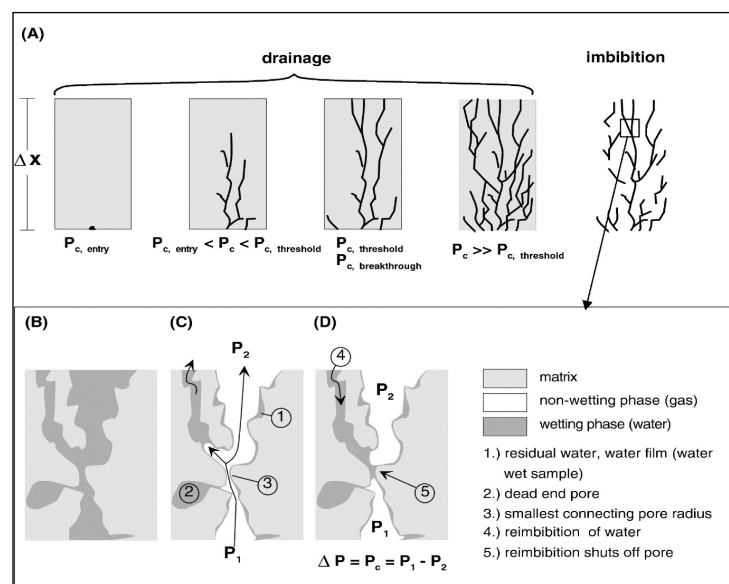


Figure 3: Re-imbibition process in fine-grained rocks (schematic re-imbibition); (A) drainage, (B) initially water-saturated sample, (C) gas breakthrough, (D) re-imbibition (Hildenbrand et al., 2002).

### 3.1.2 Effect of Temperature on the Differential Pressure Profile of GCO<sub>2</sub>-Water Displacements

Figure 4 presents the impact of increasing experimental temperature on the differential pressure profile. The results demonstrate that the increase in the experimental temperature has a significant impact on the differential pressure profile. Firstly, increasing the temperature increases the frequency of the PD oscillations. At an experimental temperature of 29 °C, the differential pressure profile experienced no oscillations. However, as the temperature increased to 31 °C, the oscillations appeared for the first time. A further increase in the temperature to 33 °C caused the number of oscillations to increase by double. Secondly, the increase in the temperature prompts an increase in the magnitude of the maximum-differential pressure. The quasi-differential pressure was almost constant due to the slight impact of both capillary forces and viscous forces at the end of core flooding.

The appearance and frequency of the PD cycles with increasing temperature have three potential explanations. The first potential reason behind the onset of the oscillations and their frequency is the increase in the capillary forces despite the slight increase in viscous forces under these conditions. The increase in temperature leads to an increase in the CO<sub>2</sub>-water IFT ([Iglauer et al., 2012](#)) with a reduction in the contact angle ([Yang et al., 2007](#)) due to the decline in the CO<sub>2</sub> solubility ([Bennion and Bachu; Yang et al., 2007](#)) as well as a slight increase in CO<sub>2</sub> viscosity, and a slight increase in CO<sub>2</sub> injection rate inside the core sample due to expansion effect. For illustration, as the experimental temperature increased from 29 to 31 °C, CO<sub>2</sub>-water IFT increases from to 42.9 to 44.42 mN/m, CO<sub>2</sub> viscosity increases very slightly from 16.72 to  $16.755 \times [10^{-6}(\text{Pa}\cdot\text{s})]$  and CO<sub>2</sub> injection inside the core sample increased from around 0.45 to 0.46 ml/min. However, a further increase in the experimental temperature to 33 °C caused the CO<sub>2</sub>-water IFT to decrease to 34.1 mN/m ([Bachu and Bennion, 2008a](#)), CO<sub>2</sub> viscosity to increase to  $16.805 \times [10^{-6}(\text{Pa}\cdot\text{s})]$  and CO<sub>2</sub> injection to increase to 0.466 ml/min.

The second possible reason might be related to the fluctuating behaviour in the CO<sub>2</sub>-water IFT when the experimental temperature is around the critical point ([Bennion and Bachu](#)), as shown in Figure 5. The third potential reason is that the PD oscillations might occur because of increasing

temperature which results in a quicker increase in the movement of CO<sub>2</sub> molecules. This is because each individual molecule has more energy as it becomes hotter, according to the Kinetic molecular theory ([Physics, 2017](#)). A high energetic CO<sub>2</sub> molecule might open the closed flow path, due to the increase in capillary forces, quicker.

The results indicate that for the sandstone core sample (from the Guillemot A field, North Sea) used in the experiment and under the aforementioned experimental conditions, the onset temperature point of the oscillations is around 31 °C. The characteristics of the sandstone sample, e.g. pore size distribution, play a key role in the onset of the PD oscillations phenomena as they have a direct influence on the magnitude of the capillary forces as illustrated by Young-Laplace law (Eq.2).

The data also reveals that as the experimental temperature increased from 29 to 31 °C, the maximum-differential pressure increased by around 12.5% (from 0.72 to 0.81 bar) and the corresponding time dropped by around 9.1% (from 12.1 to 11 min). However, increasing the temperature from 31 to 33 °C caused the differential pressure to decline slightly by 1.23% (from 0.82 to 0.81 bar) and the corresponding time dropped by 30% (from 11 to 7.7 min). The increase and decrease in the maximum-differential pressure can be related mainly to the increase or decrease in the capillary forces due to CO<sub>2</sub>-water IFT, as stated above. [The highest reduction in the corresponding time occurred as the temperature increased to 33 °C. This can be related to the highest reduction in the CO<sub>2</sub>-water IFT \(Bennion and Bachu\), as shown in Figure 5.](#)



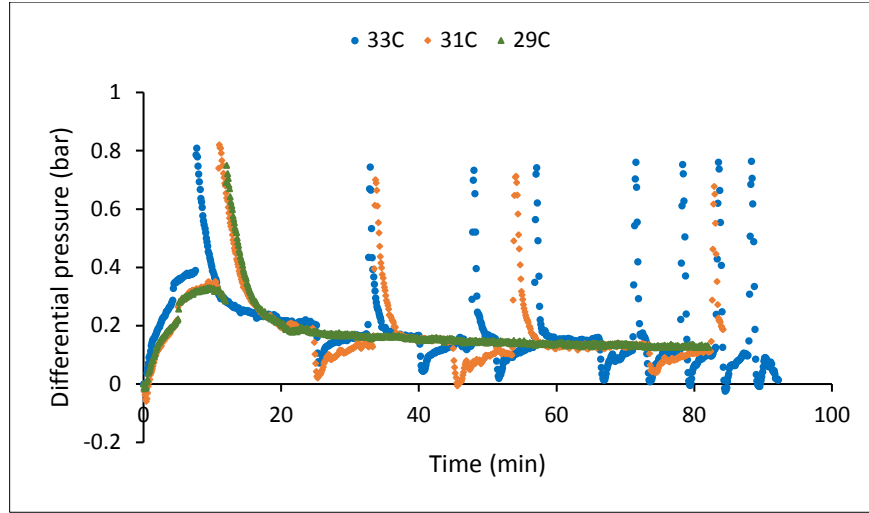


Figure 4: Effect of temperature on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 50 bar and 0.4 ml/min.

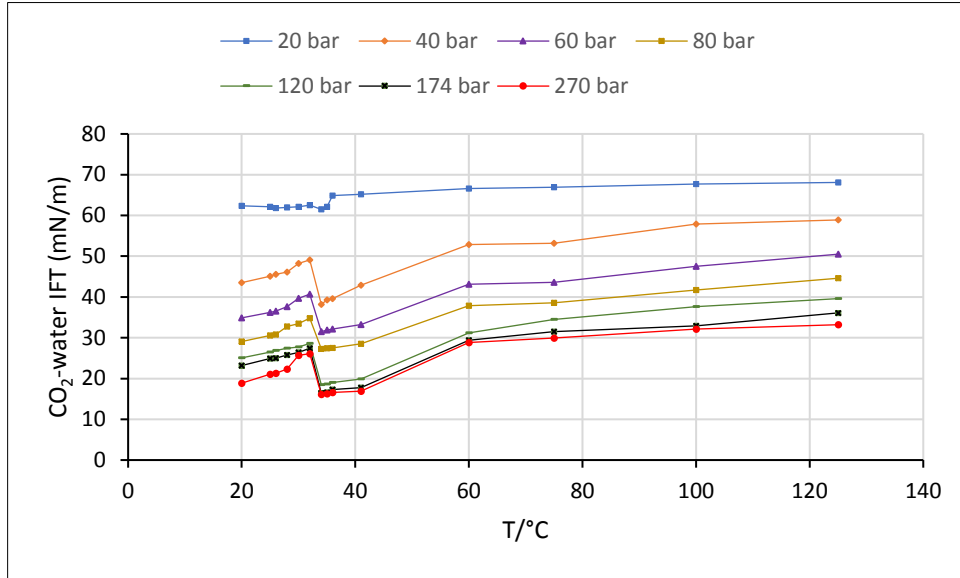


Figure 5: Interfacial tension for CO<sub>2</sub>-Pure Water Systems adopted from ([Bachu and Bennion, 2008b](#)).

To further investigate the effect of the temperature on the differential pressure profile, and especially on the PD oscillations, additional GCO<sub>2</sub>-water displacement experiments were conducted under a high-pressure of 70 bar and higher temperature conditions.

The data from Figure 6 shows that increasing the experimental temperature by 12 degrees (from 33 to 45 °C) at a high-pressure caused no further increase in the rate of the PD oscillations. Yet, it instigated a very slight increase in the maximum and quasi-differential pressures with a small reduction in the corresponding time. The maximum differential pressure increased by only 4.2% (from 0.854 to

0.89 bar) and the quasi-differential pressure by 4.81% (from 0.208 to 0.218 bar). The corresponding time declined by around 17% (from 1.8 to 1.5 min).

The data showed no further increase in the PD oscillations occurred when there are no fluctuations in the IFT as the temperature increased from 33 to 45 °C, as shown in Figure 5. This suggests that the IFT fluctuations might have highly influenced the frequency of PD oscillations.

The increase in the maximum and quasi-differential pressures can be related to the increase in the capillary forces (because of the increasing CO<sub>2</sub>-water interfacial tension and the reducing contact angle (Yang et al., 2007)), and the slight increase in the viscous forces (because of the increasing injection rate). The magnitude of the viscous forces might have slightly declined because of the slight reduction in CO<sub>2</sub> viscosity with increasing temperature. For illustration, as the experimental temperature increased from 33 to 45 °C, the CO<sub>2</sub>-water IFT increases from around 29.15 to around 33.4 mN/m (Bennion and Bachu), and the CO<sub>2</sub> injection rate inside the core sample increased from 1.315 to 1.748 ml/min but the viscosity decreases from 20.743 to 19.05 × [10<sup>-6</sup>(Pa·s)].

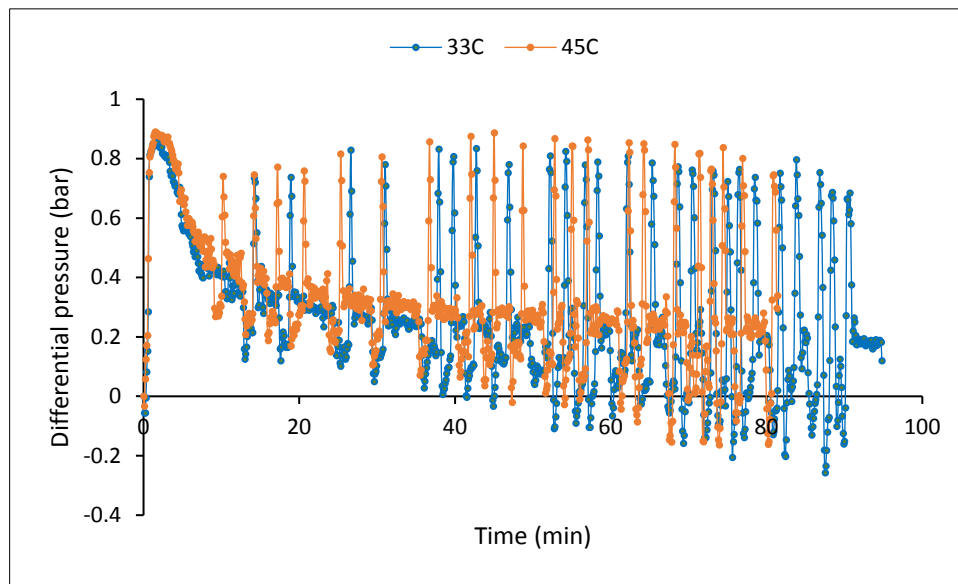


Figure 6: Effect of temperature on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 70 bar and 0.4 ml/min.

### 3.1.3 *Effect of CO<sub>2</sub> Injection Rate on the Differential Pressure Profile of GCO<sub>2</sub>-Water Core Floodings*

Figure 7, Figure 8 and Figure 9 show the impact of increasing CO<sub>2</sub> injection rate on the differential pressure profile. For Figure 8, the experiments conducted at higher injection rate (2 ml/min) lasted shorter than those conducted at lower injection rate (1 ml/min) to explore the impact of injection volumes on the displacement efficiency. The results reveal that increasing the injection rate has a significant impact on the differential pressure profile, mainly at early stages of core flooding. The data reveal a number of important observations (A-E).

A) The data show that the higher the injection rate, the higher the maximum differential pressure is. However, increasing the injection rate caused a slight increase in the quasi-differential pressure; the corresponding time decreased at low injection rates and increased at high injection rates. For illustration, as the CO<sub>2</sub> injection rate increased from 0.1 to 0.2 ml/min, the maximum-differential pressure increased by 33.54% (from 0.161 to 0.215 bar), and the quasi-differential pressure by 5.88% (from 0.068 to 0.072 bar) while the corresponding time reduced by almost half (from 13.5 to 6.5 min). However, as the CO<sub>2</sub> injection rate increased from 1 to 2 ml/min, the maximum-differential pressure increased by around 44% (from 0.833 to 1.201 bar), the quasi-differential pressure increased by around 15% (from 0.254 to 0.291 bar), and the corresponding time increased by 12% (from 3.3 to 3.7 min). The increase in the corresponding time at high injection rates despite the increase in the CO<sub>2</sub> injection rate can be related to the high increase in the magnitude of the maximum-differential pressure as well as the low-density nature of the gaseous CO<sub>2</sub>. Since the injected gaseous CO<sub>2</sub> was at low pressure (40 bar), it needed a longer time to reach the higher maximum-differential pressure of 1.201 bar during the 2 ml/min-displacement.

B) The data from Figure 7 and Figure 8 reveals that as the injection rate increased by tenfold (from 0.1 to 1 ml/min, and from 0.2 to 2 ml/min), the quasi-differential pressure increased by only around fourfold, (from 0.068 to 0.254 bar, and from 0.072 to 0.291 bar). This might be related to a potential

increase in the relative permeability with increasing injection rate ([Akbarabadi and Piri; Chang et al., 2013](#)) that leads to a reduction in the viscous pressure drop.

C) The data previously shown in [Figure 2](#) reveals that the differential pressure profile of the 40 bar-experiments is characterized by PD oscillations at 0.4 ml/min CO<sub>2</sub> injection rate. Surprisingly, the data from [Figure 7](#) and [Figure 8](#) reveal no PD oscillations at lower and higher CO<sub>2</sub> injection rates. The disappearance of the PD oscillations at higher injection rates (e.g. 1-2 ml/min) can be related to the high increase in the pressure drop due to viscous forces. Thus, the viscous forces impeded the capillary forces, which are responsible for the observed PD oscillations phenomenon ([Nutt, 1982](#)). On the other hand, at lower CO<sub>2</sub> injection rates (e.g. 0.1 to 0.2 ml/min), CO<sub>2</sub> might flow through preferential inlet and outlet pores ([Gunde et al., 2010](#)) that are characterized by low resistance to flow and by less capillary forces. Consequently, CO<sub>2</sub> does not need to pass through the smallest channels that are characterized by higher resistance to CO<sub>2</sub> flow and higher capillary forces, hence avoiding the impact of the capillary forces that cause the oscillations.

D) To look in detail at the unexpected results regarding the appearance and disappearance of the PD oscillations and the impact of CO<sub>2</sub> injection rate on the differential pressure profile, further experiments were conducted at 40 bar [and](#) over a more detailed range of injection rates, as shown in [Figure 9](#). It should be noted that the 0.4 ml/min GCO<sub>2</sub>-water displacement is repeated to make sure that the observations were not an experimental error.

The results from [Figure 9](#) show clearly that the PD oscillations occurred only at 0.4 ml/min for the experiments conducted at a low pressure of 40 bar. Overall, the data confirm that the increase in the injection rate produces an increase in the maximum-differential pressure and a reduction in its corresponding time for this range of injection rates. The quasi-differential pressure reduced slightly due to the potential increase in the relative permeabilities ([Akbarabadi and Piri; Chang et al., 2013](#)).

The data from [Figure 9](#) can be divided into two groups. The first group includes the experiments conducted at a CO<sub>2</sub> injection rate of 0.3 and 0.4 ml/min while the second group involves the experiments

performed at 0.5 and 0.6 ml/min. As the CO<sub>2</sub> injection rate increased for the first lower injection rate group, the maximum-differential pressure was almost constant at around 0.76 bar, but the corresponding time reduced by 25% (from around 20 to 15 min). The second higher injection rate group is characterized by a constant maximum-differential pressure of 0.938 bar and a constant corresponding time of 6.5 min. Thus, the data reveals that shifting the CO<sub>2</sub> injection rate to the second group caused the maximum-differential pressure to increase by 23.42% and the corresponding time to reduce by around 57%. The increase in the maximum-differential pressure associated with shifting the CO<sub>2</sub> injection rate might be related to the properties of the core sample. It might have occurred because as the injection rate increased from the first to the second group, the maximum-differential pressure had to further increase to open new preferential flow paths for the injected CO<sub>2</sub> (Gunde et al., 2010). The nearly constant maximum-differential pressure for each group might indicate a minimal impact for the viscous forces on the differential pressure at low pressures. It indicates also that the expected increase in the maximum-differential pressure due to increasing injection rate is reduced by the potential increase in the relative permeability due to the increasing injection rate (Akbarabadi and Piri; Chang et al., 2013).

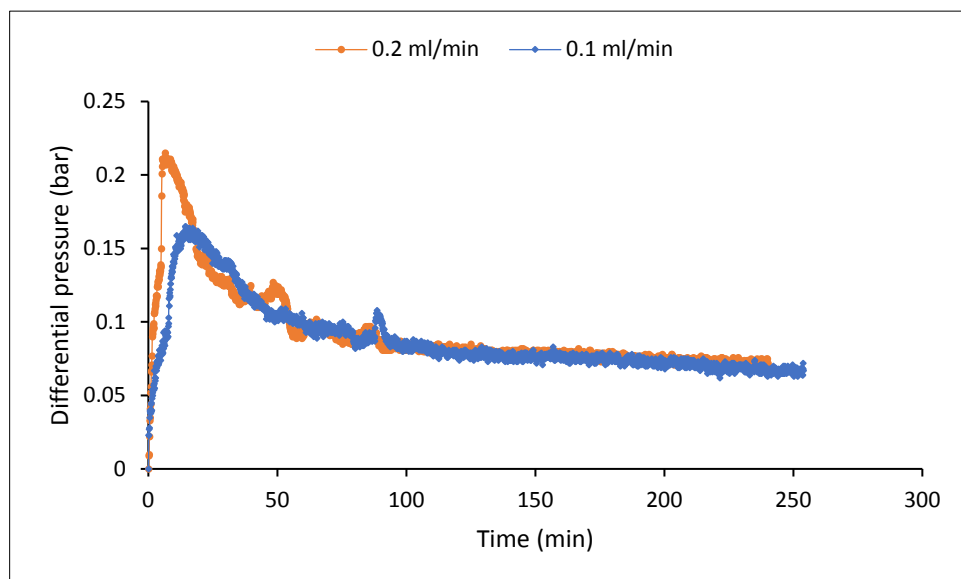


Figure 7: Effect of CO<sub>2</sub> injection rate on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 40 bar and 33 °C.

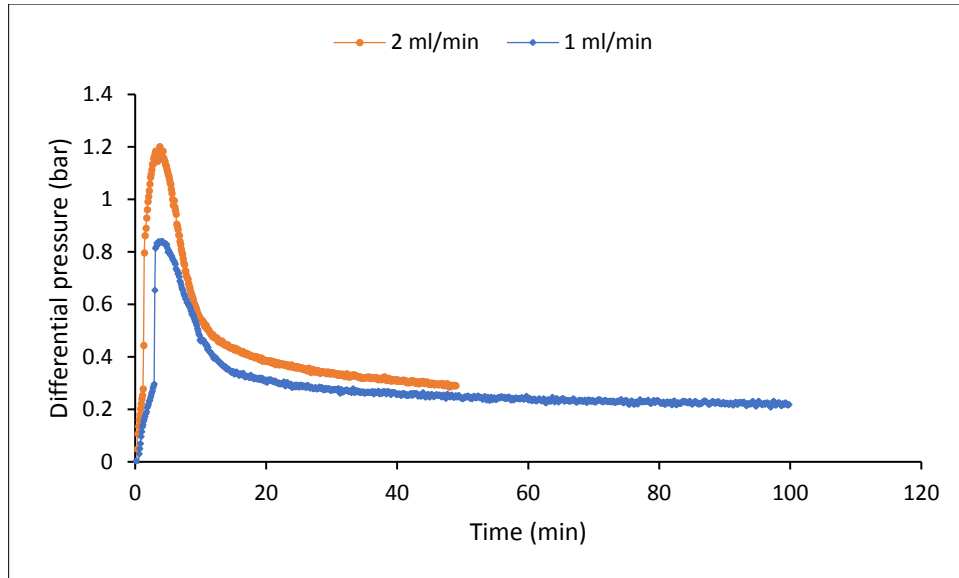


Figure 8: Effect of CO<sub>2</sub> injection rate on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 40 bar and 33 °C.

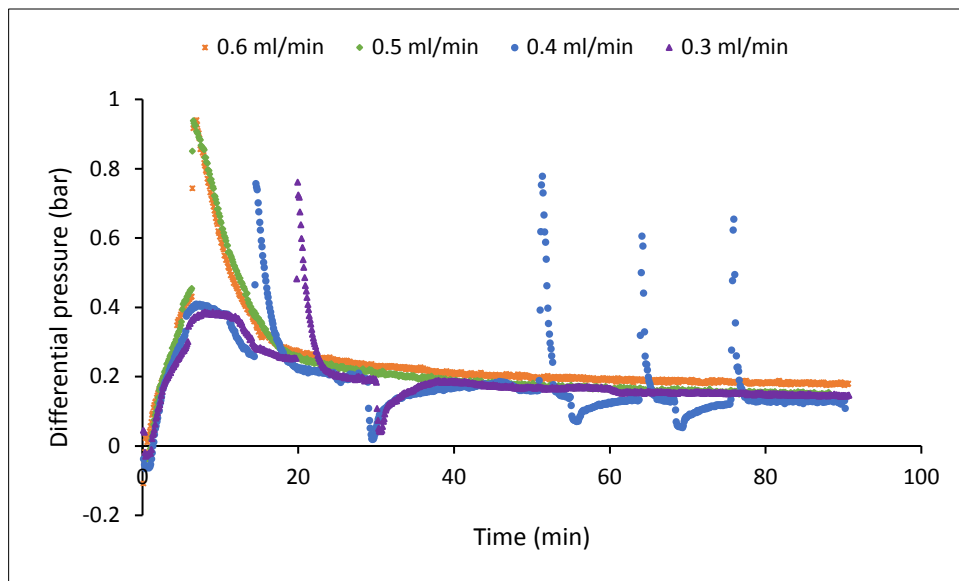


Figure 9: Effect of CO<sub>2</sub> injection rate on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 40 bar and 33 °C.

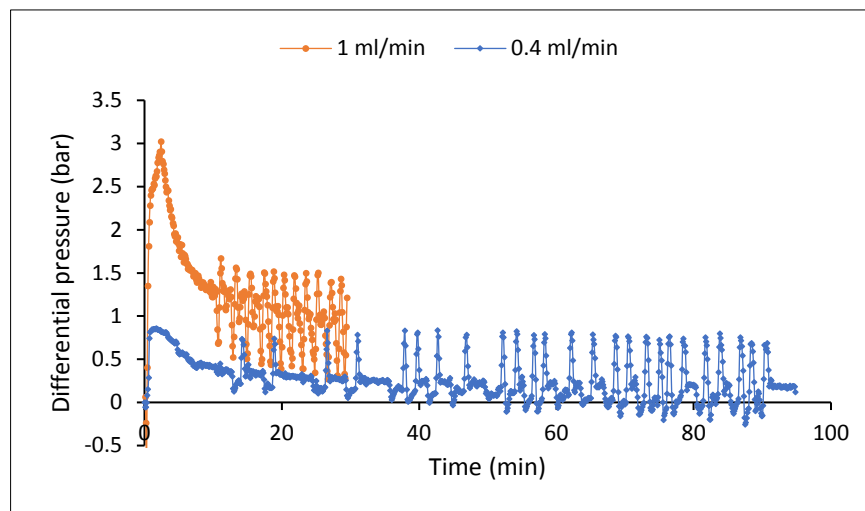
E) To further investigate the effect of CO<sub>2</sub> injection rate on the differential pressure profile and the phenomenon of the PD oscillations particularly, a second set of GCO<sub>2</sub>-water displacements have been performed at a higher pressure (70 bar). To enable a clear comparison, the data was presented in two figures: Figure 10 and Figure 11.

E.1) The data shows clearly that conducting GCO<sub>2</sub>-water displacements at higher pressure (70 bar) caused the PD oscillations to appear over a wider range of CO<sub>2</sub> injection rates (from 0.2 to 1 ml/min). It reveals also that the change in the maximum and quasi-differential pressures, corresponding time and PD oscillations depend on the range of the injection rate; the highest change occurred as the injection rate increased from 0.4 to 1 ml/min. For illustration, as the CO<sub>2</sub> injection rate increased from 0.4 to 1 ml/min, the maximum-differential pressure increased considerably by around 258% (from 0.845 to 3.024 bar) and the quasi-differential pressure increased by around 224.5% (from 0.265 to 0.86 bar). The corresponding time prolonged by 140% (from 1 to 2.4 min), despite the increase in the injection rate, due to the increase in the maximum-differential pressure. The frequency of the PD oscillations was almost constant for the last 20 min of both experiments. The increase in the maximum and quasi-differential pressures can be attributed to the increase in the viscous forces; the increase in the corresponding time can be related to the high increase in the magnitude of the maximum differential pressure.

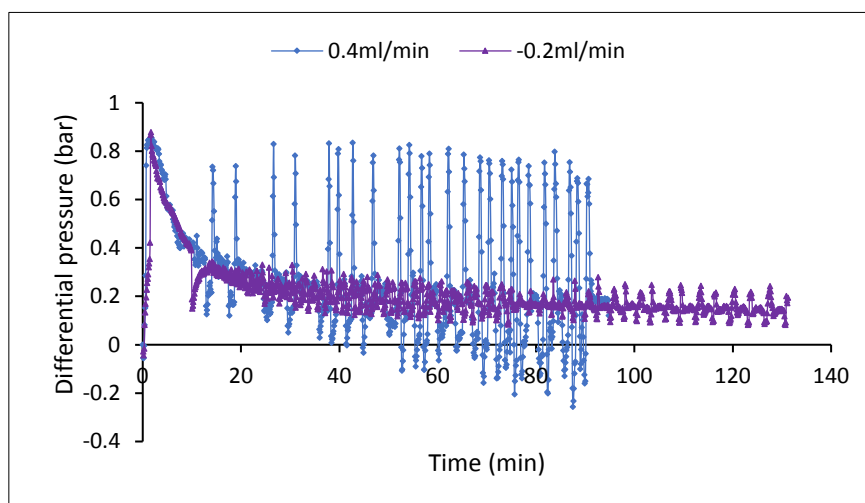
E.2) On the other hand, as the CO<sub>2</sub> injection rate increased from 0.2 to 0.4 ml/min, the maximum-differential pressure was almost constant at around 0.85 bar, the quasi-differential pressure slightly increased, the corresponding time slightly reduced, and the frequency of the PD oscillations considerably decreased but the magnitude of the PD oscillations significantly increased from around 0.25 to 0.825 bar. The nearly constant maximum-differential pressure (0.85 bar) at the low injection rates (0.2 to 0.4 ml/min)-core floodings reveals a negligible impact of the viscous forces on the differential pressure at the conditions investigated. However, the reduction in the frequency of the PD oscillations might be attributed to CO<sub>2</sub> flow through preferential flow paths ([Gunde et al., 2010](#)).

The frequency of the PD oscillations might depend to a considerable extent on the core sample properties, the change in CO<sub>2</sub> distribution due to the change in the CO<sub>2</sub> injection rate, and the operational conditions. For illustration, as the CO<sub>2</sub> injection rate increased from 0.2 to 0.4 ml/min, the CO<sub>2</sub> might have distributed over a wider range of capillaries. Consequently, as the viscous pressure drop declined because of water depletion, the CO<sub>2</sub> flow inside the smaller capillaries was blocked due

to their higher resistance to CO<sub>2</sub> flow. Later, as the pressure drop continued, the CO<sub>2</sub> flow in larger capillaries was blocked, too. Ultimately, it came to the point when all capillaries were blocked by the capillary forces ([Hildenbrand et al., 2002](#); [Nutt, 1982](#)). Thus, the increase in CO<sub>2</sub> distribution with increasing injection rate might have led to prolonging the time required for the capillary forces to block the CO<sub>2</sub> production from all opened interconnected flow paths. As a result, since the volume of the opened capillaries were larger with increasing injection rate from 0.2 to 0.4 ml/min; therefore, the frequency of the PD oscillations was reduced.



**Figure 10:** Effect of CO<sub>2</sub> injection rate on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 70 bar and 33 °C.



**Figure 11:** Effect of CO<sub>2</sub> injection rate on the differential pressure profile of GCO<sub>2</sub>-water displacements conducted at 70 bar and 33 °C.



In summary, fluid pressure, temperature and CO<sub>2</sub> injection rate exert significant influences on the differential pressure profile of the GCO<sub>2</sub>-water drainage displacements. The differential pressure profile at all fluid pressures, temperatures and injection rates is characterized by a high initial increase immediately followed by a steep rapid pressure reduction and then by a quasi-pressure drop.

The differential pressure is controlled by the interplay of both capillary and viscous forces. The increase in capillary forces leads to the appearance of the PD oscillations (the onset points) while the increase in viscous forces causes their impedance.

There are multiple cycles of these oscillations and the occurrence and frequency of these oscillations vary with fluid pressure, temperature and injection rate. The frequency of these oscillating cycles increases as fluid pressure and fluid temperature increase but vary with injection rate and seem to be fluid pressure dependent. These oscillations occurred only at 0.4 ml/min at low pressures (i.e. 40 bar), but they appeared over a wider range of injection rates at higher pressures (i.e. 70 bar). The maximum-differential pressure reached during each cycle increases with increasing fluid pressure, temperature and injection rate.

### ***3.2 Effect of Fluid Pressure, Temperature, and Injection Rate on Irreducible Water Saturation and Endpoint Effective and Relative Permeabilities of CO<sub>2</sub>***

The effective and relative permeabilities of CO<sub>2</sub> are significantly important to the determination of the efficiency and integrity of CO<sub>2</sub> sequestration in subsurface formations ([Busch and Müller, 2011](#); [Rathnaweera et al., 2015](#)). At the end of the flooding experiment, the volume of the water produced was measured, and the irreducible water saturation was calculated. Then, the core sample was weighed to confirm the irreducible water saturation calculations. To calculate the endpoint effective (relative) CO<sub>2</sub> permeability using Darcy's law, the average quasi-differential pressure and the average CO<sub>2</sub> outflow rate of the last period were used ([Akbarabadi and Piri; Chang et al., 2013](#)). The CO<sub>2</sub> viscosity at the experimental pressure and temperature was calculated using the Peace software website ([Peace software, 2017](#)).

The results from [Table 1](#) shows that both endpoint CO<sub>2</sub> relative permeability ( $K_{rCO_2}^{max}$ ) ([Armstrong et al., 2017](#)) and irreducible water saturation ( $S_{wr}$ ) are dependent on the experimental conditions at which they are measured. The  $S_{wr}$  was in the range of around 0.38-0.45 while the  $K_{rCO_2}^{max}$  was less than 0.25. Busch and Müller obtained a low relative permeability for CO<sub>2</sub>, too ([Busch and Müller, 2011](#)). Such

low relative permeability would tend to decrease injectivity while increasing displacements efficiency (Levine et al., 2011).

The results from Table 1 reveal that in general the increase in fluid pressure, temperature, and injection rate lead to an increase in the  $K_{rCO_2}^{max}$  and a decline in the  $S_{irr}$ . In case of increasing fluid pressure and temperature, the high increase in the  $K_{rCO_2}$  can be attributed mainly to the high increase in the injection rate inside the core sample due to the high impact of gas expansion (Rostami et al., 2010; Skauge et al.). This increase in volumetric CO<sub>2</sub> injection rate might result in forcing the CO<sub>2</sub> to flow through a wider range of the core sample pores.

The displacements efficiency is controlled by many factors that include relative permeability, wetting conditions, viscous fingering, gravity segregation, channelling, the amount of crossflow/mass transfer (Chukwudeme and Hamouda, 2009), mobility ratio, and capillary number (Kazemifar et al., 2015). The capillary number ( $Ca$ ) refers to the ratio of the viscous forces to capillary forces (Lenormand et al., 1988). The mobility ratio ( $M$ ) refers to the ratio of the displaced to the displacing phase viscosities. Increasing the contrast between the viscosity of the displacing and displaced fluid leads to a higher  $M$  which will result in a more unstable configuration front. The following formulas are used to define them:

$$Ca = \frac{\mu_2 V_2}{\sigma \cos \theta} \quad (5)$$

$$M = \frac{\mu_2}{\mu_1} \quad (6)$$

where  $\mu$  is the dynamic viscosity,  $\sigma$  the interfacial tension between the displaced and the displacing phases, 1 the subscript of the displaced phase, 2 the subscript of the displacing phase,  $\theta$  the contact angle between the two fluids and the surface, and  $V_2$  the bulk velocity of the displacing fluid. The flowing equation is used to define the bulk velocity.

$$V_{bulk} = \frac{Q}{A \phi} \quad (7)$$

where  $Q$  is the volumetric injection rate,  $A$  the area of the frontal face of the core sample, and  $\phi$  the core sample porosity ([Kazemifar et al., 2015](#)). Based on the magnitudes of the  $Ca$  and the  $M$ , three different regimes can be defined ([Kazemifar et al., 2015](#)). For the GCO<sub>2</sub>-water displacement investigated both  $Ca$  and  $M$  are small, which suggest a capillary fingering regime.

The reduction observed in the  $S_{wr}$  can be attributed mainly to the increase in the  $Ca$  and the reduction in the  $M$ . This is because the  $Ca$  and  $M$  are the most influential dimensionless parameters that govern GCO<sub>2</sub>-water core flooding displacement ([Kazemifar et al., 2015](#)). As the  $Ca$  increases, the impact of the capillary forces compared to viscous forces decreases. The balance between the viscous forces and capillary forces governs the pore scale drainage displacements ([Heaviside and Black, 1983](#)). The capillary forces are responsible for the trapping of the injected CO<sub>2</sub> ([Akbarabadi and Piri, 2013](#); [Bachu and Shaw, 2003](#)). Thus, decreasing the capillary forces (e.g. due to the reduction in the interfacial tension) will lower the  $S_{wr}$  (i.e. enhance the fluid displacements) ([Ahmadi et al., 2015](#)). On the other hand, reducing  $M$  will result in a more uniform displacement of water by CO<sub>2</sub> ([Bennion and Bachu, 2006](#)), which can result in reducing the  $S_{wr}$ . The data from [Table 1](#) show that the increase in the  $Ca$  and the reduction in the  $M$  can lead to a reduction in the  $S_{wr}$  even when the change in both  $Ca$  and  $M$  is small. Ding and Kantzas observed that the critical  $Ca$  for the gas-water system is 2E-8 ([Ding and Kantzas, 2007](#)).

The results from [Table 1](#) reveal that increasing the fluid pressure from 40 to 70 bar at 33 °C and 0.4 ml/min caused the  $K_{rCO_2}^{max}$  to increase by around 0.099 and the  $S_{wr}$  to decrease by around 0.047. The largest increase in the  $K_{rCO_2}^{max}$  and the highest reduction in the  $S_{wr}$  occurred as the fluid pressure increased from low-fluid pressure displacements (40 and 50 bar) to high-fluid pressure displacements (70 bar). The observed trend of the  $K_{rCO_2}^{max}$  and  $S_{wr}$  are in agreement with the findings of Liu et al. and Bennion and Bachu ([Bennion and Bachu, 2006](#); [Liu et al.](#)). Liu et al also observed an increase in the  $K_{rCO_2}$  with increasing pressure ([Liu et al.](#)). Bennion and Bachu observed an increase in the  $K_{rCO_2}$  and increase in the maximum endpoint CO<sub>2</sub> saturation (i.e. decrease in  $S_{wr}$ ) with increasing pressure; they attributed

that to the reduction in IFT with increasing pressure ([Bennion and Bachu, 2006](#)). The observed trend of the  $K_{rCO_2}^{max}$  and  $S_{wr}$  can also be associated with the relatively high increase in the  $Ca$  and the high reduction in the  $M$ .

The results from [Table 1](#) reveal that increasing temperature led to an increase in the  $K_{rCO_2}^{max}$ . On the other hand, increasing temperature caused a reduction in the  $S_{wr}$  for the displacements conducted at high-fluid pressure (70 bar) and over a high temperature increase (33-45 °C). Nonetheless, for the experiments conducted at low-fluid pressure (50 bar) and over a small temperature increase (29-33 °C), the trend of the  $S_{wr}$  depends on the magnitude of the experimental temperature. For the high-fluid pressure displacements, when the temperature increased from 33 to 45 °C at 70 bar, the  $K_{rCO_2}^{max}$  increased by around 0.035 and the  $S_{wr}$  decreased by around 0.02. The reduction in the  $S_{wr}$  for the 70 bar displacements can be attributed also to the high increase in the  $Ca$  and the high reduction in the  $M$ . For the low-fluid pressure displacements, as the temperature increased slightly from 29 to 33 °C at 50 bar, the  $K_{rCO_2}^{max}$  increased by around 0.016. Nevertheless, the  $S_{wr}$  value was between around 0.40 and 0.41. The  $S_{wr}$  saturation slightly increased by around 0.01 as the temperature increased from 29 to 31 °C, and then slightly decreased by about 0.005 as the temperature increased from 31 to 33 °C. The slight increase in the  $S_{wr}$  might be related to the slight reduction in the  $Ca$  as well as the impact of the capillary forces, which can be seen through the appearance of the PD oscillations when the temperature increased to 31 °C, see Section 3.2 for more information; the PD oscillations might result in hindering water production to a slight extent. On the other hand, the slight reduction in the  $S_{wr}$ , when the temperature further increased to 33 °C, can be associated with the relatively high increase in the  $Ca$  as well as the slight reduction in the  $M$ .

Overall, the results from [Table 1](#) shows that the increase in the  $CO_2$  injection rate caused an increase in the  $K_{rCO_2}^{max}$  and a reduction in the  $S_{wr}$ . Increasing the injection rate from 0.1 to 2 ml/min at 40 bar and 33 °C resulted in an increase in the  $K_{rCO_2}^{max}$  by around 0.157 and a reduction in the  $S_{wr}$  by around 0.05. These findings agree with those in Chang et al. and Akbarabadi and Piri ([Akbarabadi and Piri; Chang](#)

et al., 2013). However, for the core flooding at 0.4 ml/min or less, the  $S_{wr}$  trend is not clear. Moreover, the  $K_{rCO_2}^{max}$  of the experiments conducted at 40 bar-0.2 ml-33 °C does not fit linearly in the trend. Increasing the injection rate from 0.6 to 1 ml/min resulted in the highest reduction in the  $S_{wr}$ . This can be corresponded to the high increase in the  $Ca$  from around 7.9 E-8 to 1.3 E-7. For the core flooding performed at 70 bar and 33 °C, increasing the injection rate from 0.2 to 1 ml/min caused a very slight reduction in the  $S_{wr}$  by 0.0077. However, the  $K_{rCO_2}^{max}$  increased substantially as the injection rate increased from 0.2 to 0.4 ml/min. Nevertheless, as the injection rate increased to 1 ml/min, a significant reduction in the  $K_{rCO_2}^{max}$  happened again, the reason is not clear. The very slight reduction in the  $S_{wr}$  might be because only a slight increase occurred in the  $Ca$  and that  $M$  was constant.

Table 1: Effect of fluid pressure, temperature, and injection rate on endpoint effective and relative permeabilities of gaseous CO<sub>2</sub> and irreducible water saturation

Parameter	Experiment	$K_{rCO_2}$	$K_{rCO_2}$	$S_{wr}$	$M$	$Ca$
Fluid Pressure Effect	40 bar-0.4 ml/min-33 °C	1.768	0.113	0.4244	46.26	5.265E-08
	50 bar-0.4 ml/min-33 °C	1.987	0.127	0.4089	44.56	6.250E-08
	70 bar-0.4 ml/min-33 °C	2.613	0.212	0.3779	36.10	2.504E-07
Temperature Effect	50 bar-0.4 ml/min-29 °C	1.507	0.096	0.4012	48.69	4.748E-08
	50 bar-0.4 ml/min-31 °C	1.738	0.111	0.4147	46.57	4.698E-08
	50 bar-0.4 ml/min-33 °C	1.987	0.127	0.4089	44.56	6.250E-08
	70 bar-0.4 ml/min-33 °C	2.613	0.212	0.3779	36.10	2.547E-07
	70 bar-0.4 ml/min-45 °C	3.675	0.247	0.3566	31.34	2.714E-07
Injection Rate Effect	40 bar-0.1 ml/min-33 °C	0.67	0.043	0.38	46.26	1.316E-08
	40 bar-0.2 ml/min-33 °C	1.265	0.081	0.446	46.26	2.632E-08
	40 bar-0.3 ml/min-33 °C	0.955	0.061	0.436	46.26	3.948E-08
	40 bar-0.4 ml/min-33 °C	1.493	0.095	0.4244	46.26	5.265E-08
	40 bar-0.5 ml/min-33 °C	1.528	0.097	0.436	46.26	6.581E-08
	40 bar-0.6 ml/min-33 °C	1.535	0.098	0.4167	46.26	7.897E-08
	40 bar-1 ml/min-33 °C	1.793	0.114	0.3837	46.26	1.316E-07

40 bar-2 ml/min-33 °C	3.13	0.20	0.391	46.26	2.632E-07
70 bar-0.2 ml/min-33 °C	2.421	0.154	0.3798	36.10	1.273E-07
70 bar-0.4 ml/min-33 °C	3.625	0.167	0.3779	36.10	2.547E-07
70 bar-1 ml/min-33 °C	1.976	0.128	0.3721	36.10	6.368E-07

#### 4. Conclusion

In this paper, the effect of fluid pressure, temperature, and CO<sub>2</sub> injection rate on gaseous CO<sub>2</sub> dynamic behaviour during its flooding of a water-saturated sandstone core sample have been investigated in detail. The results indicate that the parameters investigated have a moderate to significant influence on the differential pressure profile, endpoint CO<sub>2</sub> relative and effective permeabilities and irreducible water saturation.

For all fluid pressures, temperatures, and injection rates, the differential pressure profiles are characterized by a sharp increase, immediately followed by a steep pressure reduction, and finally, by a gradual pressure reduction. The differential pressure profiles are controlled by the interplay of both capillary and viscous forces. The capillary forces produce cyclic oscillations within the differential pressure and fluid production data; the increase in the viscous forces impede the appearance of these oscillations. The appearance and frequency of the oscillations depend on the fluid pressure, temperature, and CO<sub>2</sub> injection rates. In general, the frequency of the oscillations increased with increasing pressure and temperature. The differential pressure oscillation cycles exhibit a very interesting response to varying injection rate, they are dependent on the fluid pressure. At 40 bar, the oscillations were only observed at an injection rate of 0.4 ml/min, whereas at 70 bar the oscillations occurred at all injection rates tested (0.2, 0.4, and 1ml/min).

In general, the increase in fluid pressure, temperature, and injection rate led to an increase in the maximum and quasi-differential pressures; the extent of the increase in the differential pressure is dependent on the fluid pressure, temperature, and injection rate. Increasing the fluid pressure and temperature caused a reduction in the time required to achieve the maximum-differential pressure at

the start of the experiment, i.e. corresponding time. Whereas, increasing the injection rate caused the corresponding time to decrease at low injection rates and increase at high injection rates.

In general, the increase in fluid pressure, temperature, and injection rate led to an increase in the endpoint CO<sub>2</sub> relative permeability ( $K_{rCO_2}^{max}$ ) and a decline in the irreducible water saturation ( $S_{wr}$ ). The  $S_{wr}$  was in the range of around 0.38-0.45 while the  $K_{rCO_2}^{max}$  was less than 0.25.

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